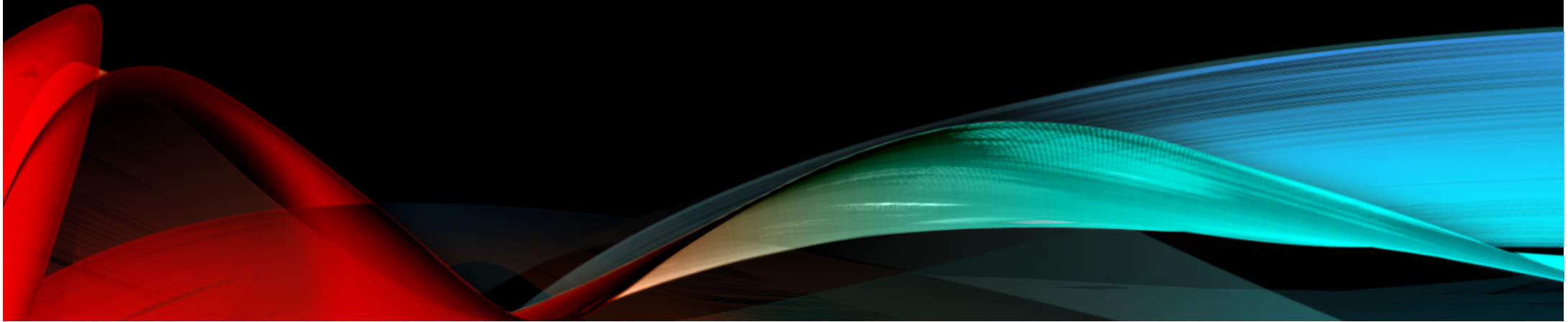
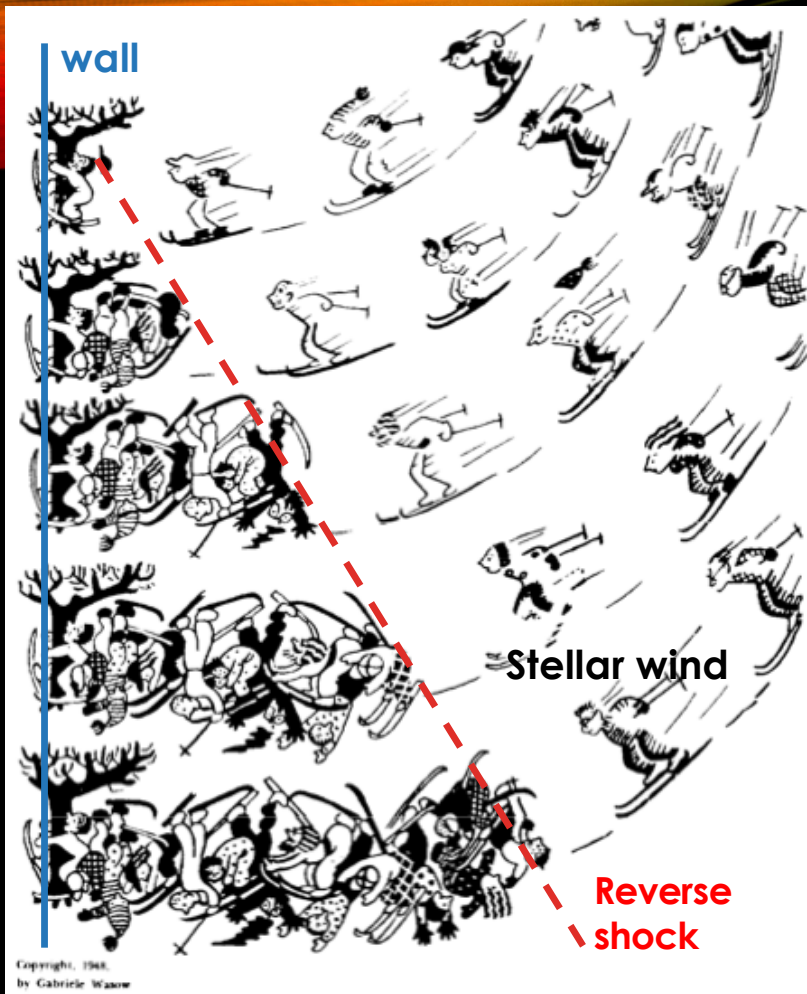


MODELING INTERSTELLAR SHOCKS

Sylvie CABRIT (Observatoire de Paris, LERMA)



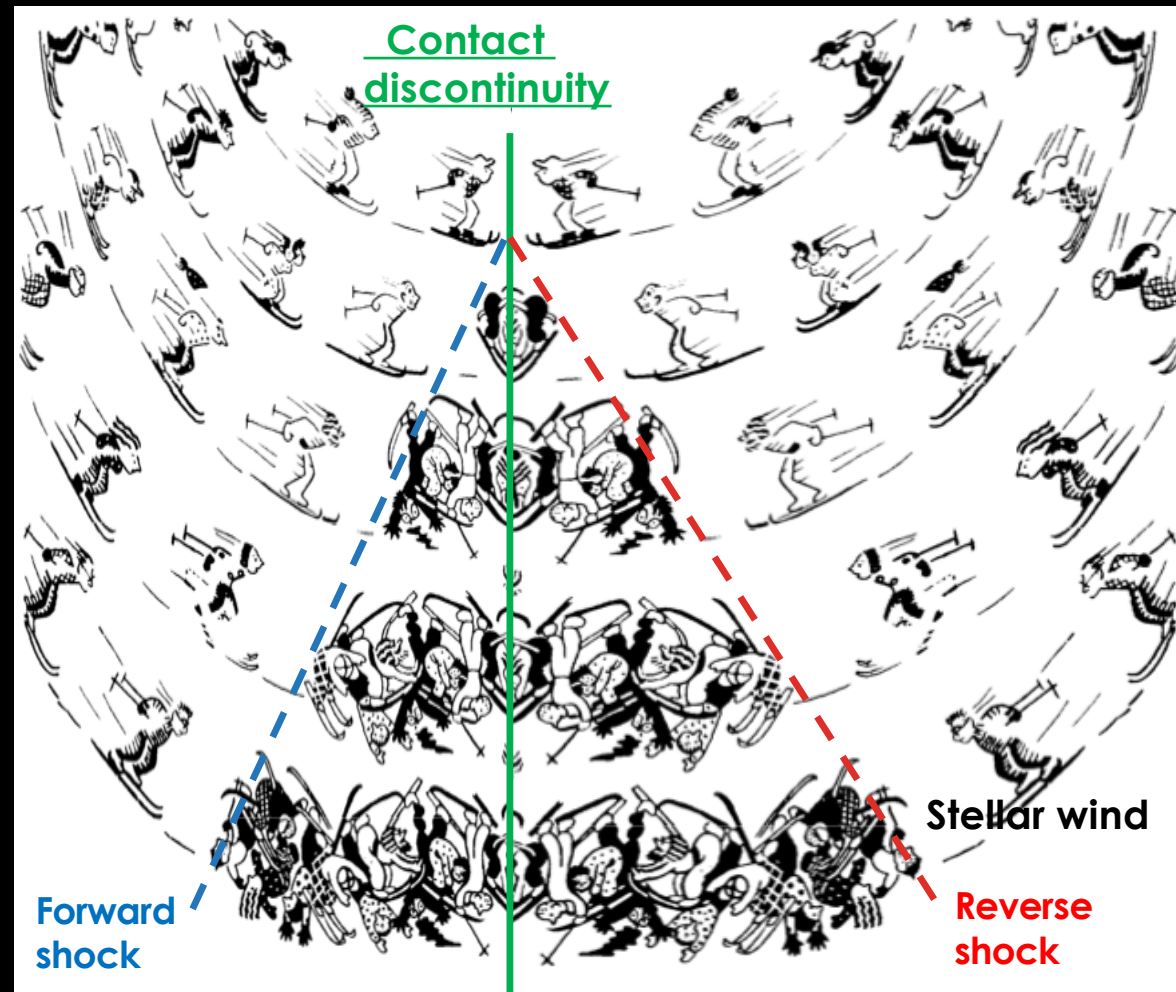


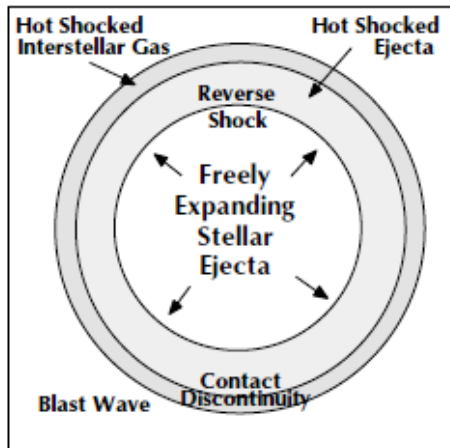
BASIC SHOCK THEORY

- Supersonic flow meets an obstacle at $V_s > \text{sound speed}$
- Sudden compression \rightarrow deceleration (mass conservation)
- Conversion of kinetic energy flux into
 - Thermal energy (heat)
 - Dissociation
 - Ionization
 - Magnetic energy (B compressed)
 - Internal energy \rightarrow **Radiation signature**

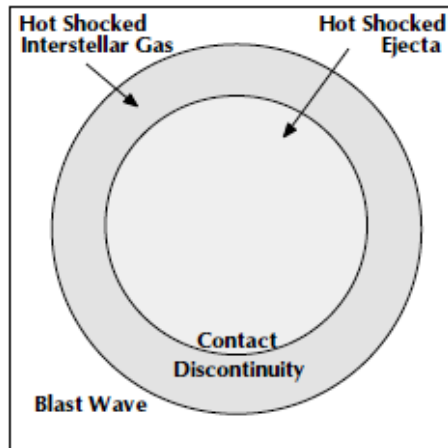
If compressible obstacle (ambient cloud, previous ejecta..)

- « Forward shock » launched into the obstacle
- « Contact discontinuity » between the two shock fronts

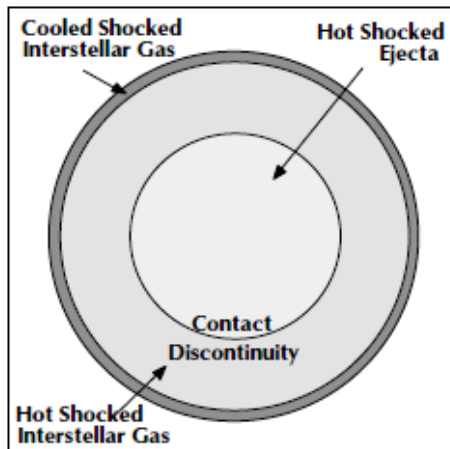




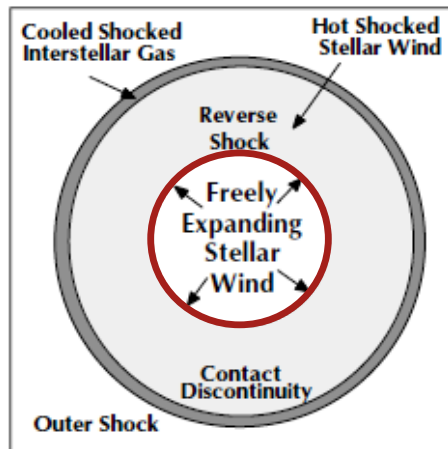
a. Supernova Remnant: pre Sedov-Taylor Phase



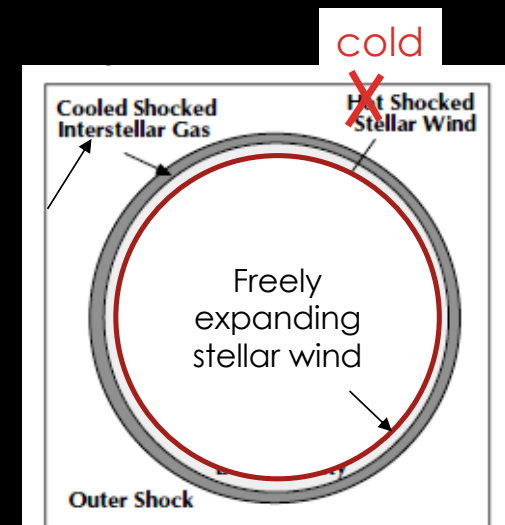
b. Supernova Remnant: Sedov-Taylor Phase



c. Supernova Remnant: post Sedov-Taylor Phase



d. Stellar Wind Bubble



e. Slow and/or dense winds (protostars, AGB stars..) : **reverse shock** **cools rapidly** → **thin shell** driven by wind ram-pressure = « snowplow »

THIN SHELL SPEED (SNOW-PLOW)

- Shell mass increase (shocked ambient gas + shocked wind)
- Shell momentum increase (from shocked wind)

• $dv_s/dt = 0 \iff$ ram pressure equilibrium between reverse and forward shock

- Solution:
- $v_s \ll v_w$ if $\rho_a \gg \rho_w$

$$\frac{d\mathcal{M}}{dt} = r^2 \rho_a v_s + r^2 \rho_w (v_w - v_s)$$

$$\frac{d\mathcal{M}v_s}{dt} = r^2 \rho_w (v_w - v_s)v_w$$

$$\rightarrow \mathcal{M} \frac{dv_s}{dt} = \frac{d\mathcal{M}v_s}{dt} - v_s \frac{d\mathcal{M}}{dt}$$

$$= r^2 \rho_w \left((v_w - v_s)^2 - \frac{\rho_a}{\rho_w} v_s^2 \right)$$

$$\rightarrow v_s \sim \frac{v_w}{1 + \sqrt{\rho_a/\rho_w}}$$

SHOCK MODELING

« Forward modeling »

- Specify driving piston + ambient medium conditions
 - compute shock front(s) structure and propagation
 - compute emitted spectrum

« Backwards modeling » (= Inversion)

- Specify emission fluxes (observed)
 - fit shock parameters
 - constrain piston properties (eg. wind mass-flux, momentum flux, power, ...)
 - constrain ambient medium properties (eg. B-field, density, ice composition,...)
 - Constrain global **shock feedback** on ISM (turbulent support, chemistry, SF regulation...)

FORWARD MODELING

Semi-Analytical

1D, steady-state

PRO:

- Very complete non-equilibrium micro-physics : ionization, chemistry, grains, ion-neutral drift...
- Fast: can compute large model grids
- Pseudo 2D = collection of 1D shocks

LIMITS:

- Ignore lateral loss of pressure in 2D
- Ignore time-dependence and instabilities

Numerical simulations

HD or MHD 1D/2D/3D

PRO:

- Exact time-dependence
- Study instabilities and turbulence
- Accurate effect of 2D curvature

LIMITS: CPU cost

- Simplified cooling and chemistry
- No large grids
- Challenging to resolve cooling length (J-shock) or treat Multi-fluid MHD (C-shock)

CONSERVATION LAWS

In the frame of shock wave

Single-fluid 1D steady-state with $B \parallel$ shock front

- **Mass conservation**

$$\rho V = \text{cste}$$

- **Magnetic Flux conservation**

$$B / \rho = \text{cste}$$

- **Momentum conservation** (ram, thermal and magnetic pressure)

$$\rho V^2 + P + B^2 / 8\pi = \text{cste}$$

- **Energy conservation**

$$V \left[\frac{1}{2} \rho V^2 + P + u + B^2 / 4\pi \right] + F_{\text{rad}} = \text{cste}$$

u = internal energy

F_{rad} is radiative cooling flux

See eg. Hollenbach & McKee 1979 for oblique B case

Rankine-Hugoniot Jump conditions

set $F_{\text{rad}} = 0$ (adiabatic front), case $B=0$

$$\rho_s = \rho_0 \frac{(\gamma + 1) \mathcal{M}^2}{(\gamma - 1) \mathcal{M}^2 + 2},$$

$$\rho_s \sim 4\rho_0 \text{ in the strong shock limit, } \gamma = 5/3, \mathcal{M} \gg 1.$$

$$\frac{T_1}{T_0} = \frac{[(\gamma - 1) \mathcal{M}_0^2 + 2][2\gamma \mathcal{M}_0^2 - (\gamma - 1)]}{(\gamma + 1) 2 \mathcal{M}_0^2}.$$

$$T_s = \frac{2(\gamma - 1)}{(\gamma + 1)^2} \left[\frac{\mu m_u v_s^2}{k} \right],$$

$\sim 3000 \text{ K @ } 10 \text{ km/s}$
 $\sim 3 \cdot 10^5 \text{ K @ } 100 \text{ km/s}$

INTEGRATED EQUATIONS

Single-fluid 1D steady-state with $B \parallel$ shock front

In the frame of the shock wave:

- **Mass conservation**

$$\rho V = \text{cste}$$

- **Magnetic Flux conservation**

$$B / \rho = \text{cste}$$

- **Momentum conservation** (ram, thermal and magnetic pressure)

$$\rho V^2 + P + B^2 / 8\pi = \text{cste}$$

- **Energy conservation**

$$V \left[\frac{1}{2} \rho V^2 + P + u + B^2 / 4\pi \right] + F_{\text{rad}} = \text{cste}$$

u = internal energy

F_{rad} is radiative cooling flux

See eg. Hollenbach & McKee 1979 for oblique B case

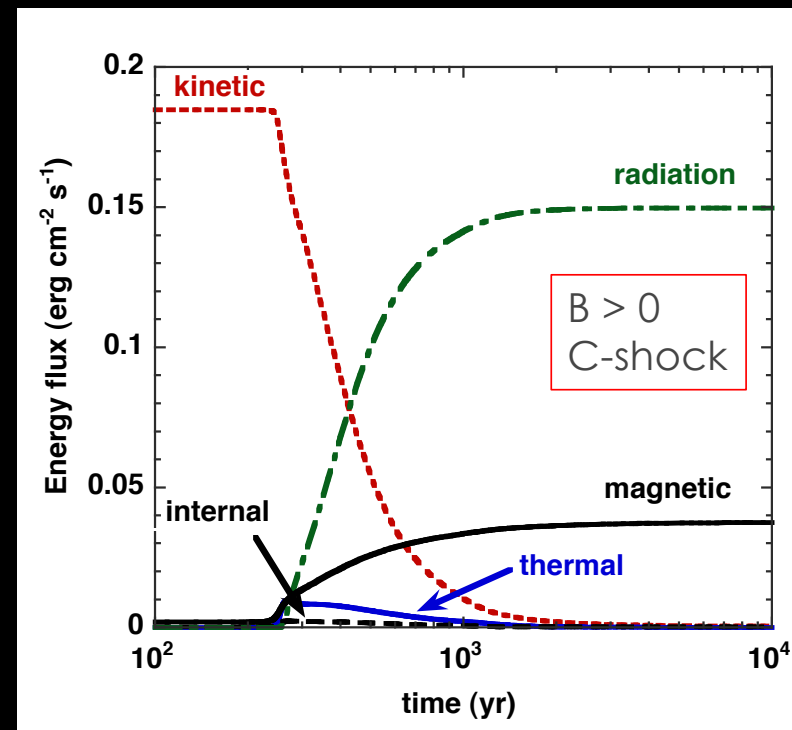
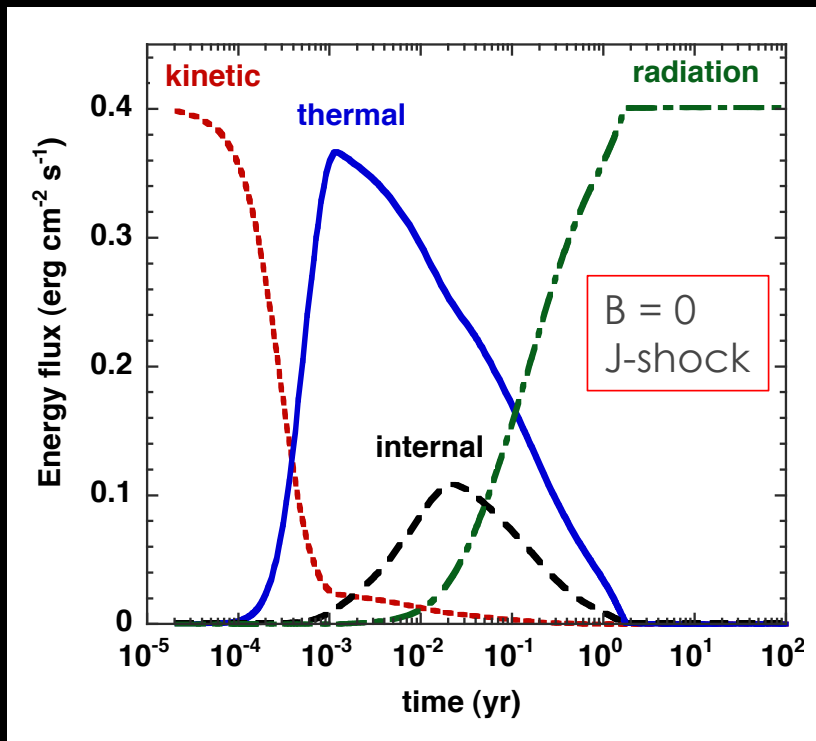
- Starting from initial conditions, solve differential equations for ρ , V , P (or T), u , F_{rad} along z
- $dF_{\text{rad}} / dz = \Lambda$ (erg/s/cm³) = cooling rate
- $P = \rho kT / \mu mH$

At equilibrium, Λ , u and μ depend only on ρ , T , elemental abundances: easy to integrate

BUT In shock cooling zone, ionization stages, molecular abundances, and internal excitation are **out of EQ** (no time to adjust)

→ Also need to solve along z for each specie abundance, ionization stage, (level populations): « chemical » source terms

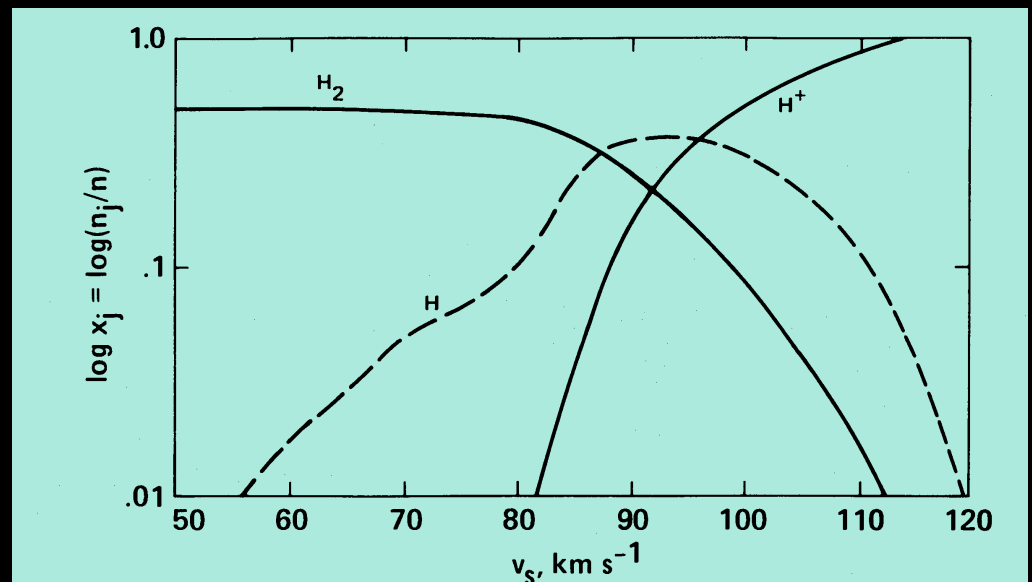
ENERGY CONVERSION IN SHOCKS



RADIATIVE PRECURSOR IN IONIZING SHOCKS

- UV photons emitted in shock front also propagate upstream in the preshock medium
→ « radiative precursor »
- $V_s < 40$ km/s: cold neutral
- $40 < V_s < 80$ km/s: warm neutral, partly dissociated
- $80 < V_s < 120$ km/s : partly ionized and fully dissociated
- 120 km/s $< V_s$: fully ionized at IEQ

Sutherland & Dopita (2017)



Hollenbach & Mc Kee (1989)

REFERENCE 1D-SHOCK MODELS

Atomic shocks, no molecules

$30 < V_s < 2000 \text{ km/s}$

- MAPPINGS V (Sutherland & Dopita 2017)

Public code available at
<https://mappings.anu.edu.au/code/>

$10 < V_s < 100 \text{ km/s}$

- Model grid in Hartigan et al. (1994, 1995)

Dissociative, with molecule reformation

$30 < V_s < 120 \text{ km/s}$

- Hollenbach & McKee 1989 (infrared line predictions)

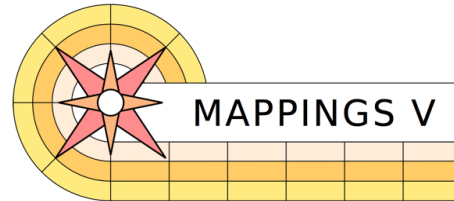
Non- (or weakly) dissociative

$V_s < 40 \text{ km/s}$

- No radiative precursor
- Paris-Durham shock code (Flower & Pineau des Forêts 1985 → Godard 2019)

<https://ism.obspm.fr/shock.html>

<https://mappings.anu.edu.au/code/>



An astrophysical plasma modelling code.

The MAPPINGS V Code Archive

<https://mappings.anu.edu.au/code>

<https://bitbucket.org/RalphSutherland/mappings>

Science and Programming:

Ralph Sutherland, Mike Dopita, David Nicholls, Brent Groves, Luc Binette, et al

Current MAPPINGS Version V 5.1.18

v5.1.18, September 2019

[Download mappings_V-5118.zip \(297.8MB\)](#) : Everything including stellar/agn atmospheres.

SHA256 digest: 9242848db55898f6bb9359f5a716ff9b52326750db5397dafc0852cf3cb6f5d22

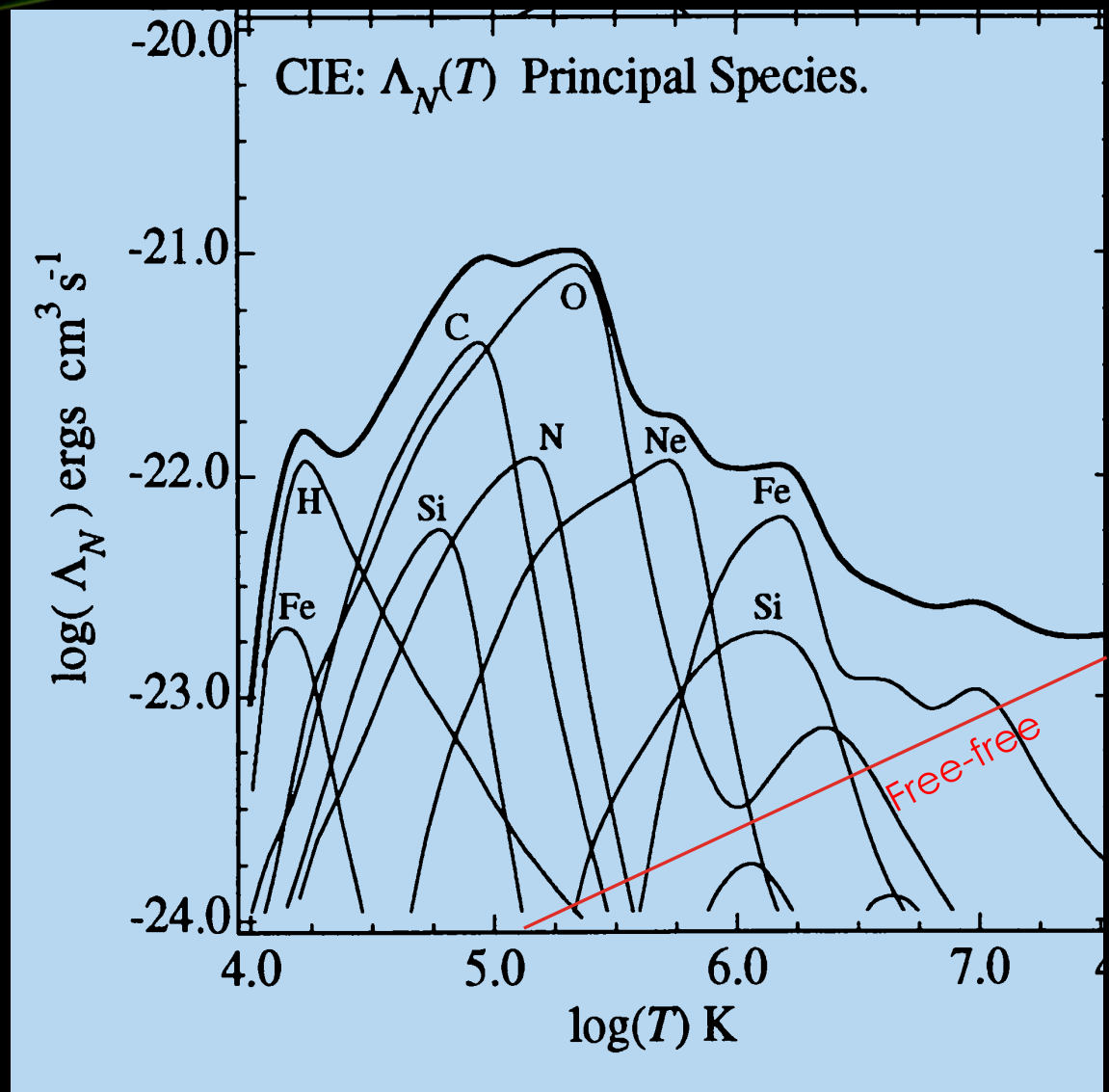
[Download M_V-5118.zip \(28.7MB\)](#) : Everything except stellar/agn atmospheres.

SHA256 digest: be429425e54baf0f23210b09a716ff9b52326750db5397dafc0852cf3cb6f5d22

MAIN COOLANTS AT $T > 10^4$ K

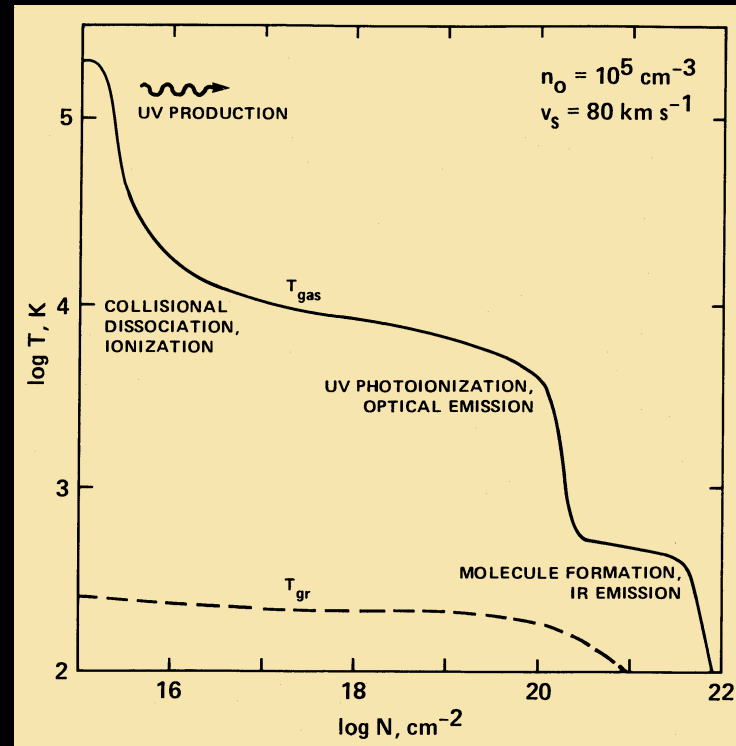
Dopita+1993 cooling curve
(MAPPINGS II)

Updated with CHIANTI database
In MAPPINGS V



DISSOCIATIVE SHOCK WITH MOLECULE REFORMATION

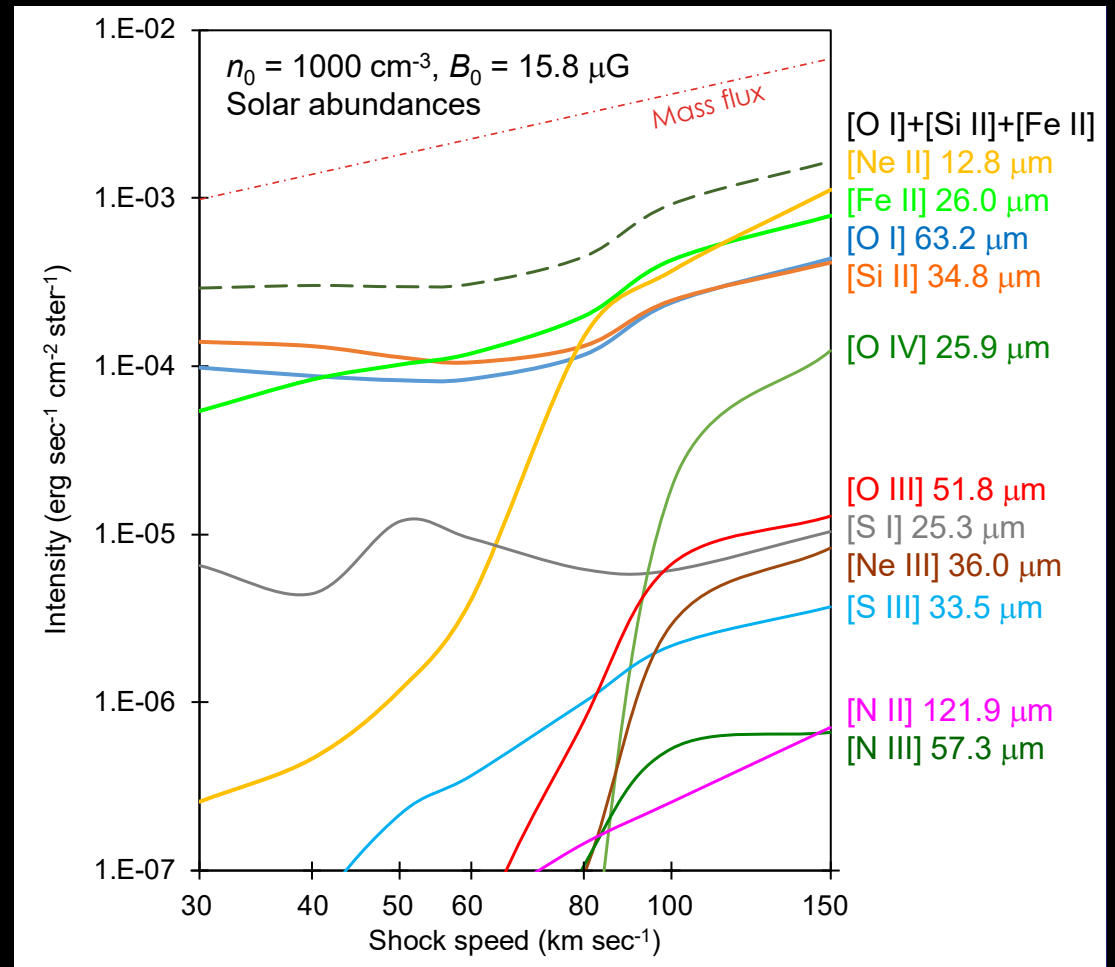
- UV photons emitted in shock front reabsorbed downstream
 - « Recombination plateau » @ 10^4 K (~ HII region)
 - Hydrogen lines and low-excitation ionic lines [O I], [Si II], [Fe II], etc..
- Molecule reformation plateau
 - @ 500 K (each H_2 molecule releases 4.48eV)
 - H_2 , CO, H_2O , (+SiO, ..)



Hollenbach & Mc Kee (1989)

SHOCK SPEED DIAGNOSTICS

MAPPINGS V calculations



[Shock code](#)
[Download](#)
[Run online](#)

Paris-Durham shock code

Download

Shock code current versions - updated in January 2019

Version	Access	Note
Shock 1.1 rev 89 (January 2019)	Download	Associated paper: Godard et al. 2019 <ul style="list-style-type: none"> • computation of C* and CJ stationary shocks • change grain treatments - adsorption / erosion widths • add thermal desorption mechanisms • update chemical network • add photoelectric effect for grain charge • update photoreaction rates with those given by Heays et al. (2017) • update the computation of secondary photon processes • add option to set up an external grid of radiation field • include radiation field spectrum and radiative transfer • integrate photoreaction using cross sections • include dust absorption properties and heat capacities • compute the grain temperature through thermal balance • add radiative pumping of H2 and CO electronic lines • compute H2 and CO self-shielding using the FGK approximation • change the way of computing H2 heating or cooling • modification of dvoid for absolute error control on log variables • possibility of sorting reactions when computing the derivatives • possibility of enforcing elements conserv for unbalanced network • add individual species velocities for S and SH • introduce an energy criterium to compute the shock size • change input / output file - plug ISM services
Shock 1.0 rev 90	Download	Associated paper: Lesaffre et al. 2013

Documentation

- [Documentation shock](#)
- [Tutorial](#)

Requirements

- Fortran 90 with lapack and blas libraries
- Python 2.7 with numpy, h5py and PyQt4

Contacts

- support.shock.ism at obspm.fr
- Sylvie.Cabrit at obspm.fr
- Antoine.Gusdorf at Ira.ens.fr
- Benjamin.Godard at phys.ens.fr

Selected references

- Flower et al., 1985, MNRAS, 216, 775
- Flower et al., 1986, MNRAS, 218, 729
- Pineau des Forêts et al., 1986, MNRAS, 220, 801
- Le Bourlot et al., 2002, MNRAS, 332, 985
- Flower & Pineau des Forêts, 2003, MNRAS, 343, 390
- Lesaffre et al., 2004, A&A, 427, 157
- Gusdorf et al., 2008, A&A, 482, 809
- Lesaffre et al., 2013, A&A, 550, A06
- Godard et al., 2019, A&A, 622, A100

SPEED LIMITS IN MAGNETIZED SHOCKS

Alfven speed in neutrals

- $V_{\text{Alfven}} = B / \sqrt{4 \pi \rho} = 1.8 \mathbf{b}$ km/s
- $M_A = V_s / V_{\text{Alfven}}$
- Maximum compression in postshock gas ($M \gg 1$)

$$C = 4 M_A^2 / (1 + \sqrt{1 + 8M_A^2})$$

→ Shock compression only if $M_A > 1$

Magnetosonic speed in charged fluid

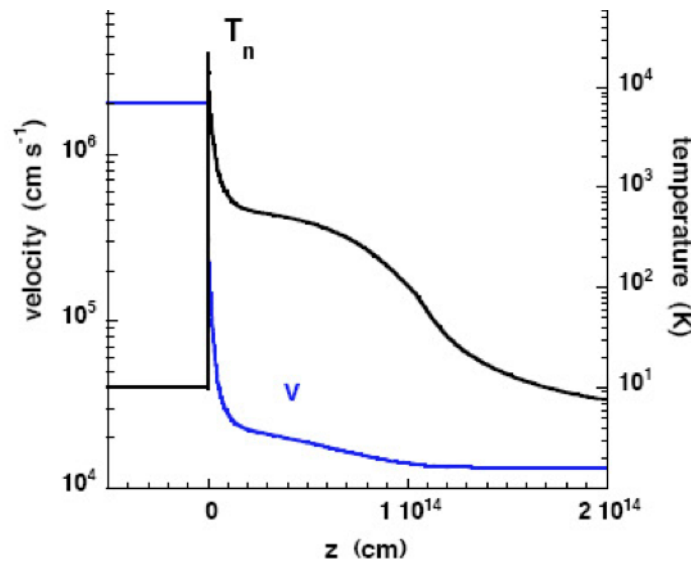
- $V_{\text{magnet}} = \sqrt{C_i^2 + B^2 / 4 \pi \rho_i}$
- If $V_A < V_s < V_{\text{magnet}}$, magnetic compression wave propagates ahead of shock in charged fluid → « magnetic precursor »
- Ions slow down by mag. pressure
- Neutrals slow down by collisions with ions → jump can disappear

J-TYPE VS C-TYPE SHOCKS

Draine (1980), Flower et al. (1985), Smith (1991), Kaufman & Neufeld (1996), Caselli et al (1997)

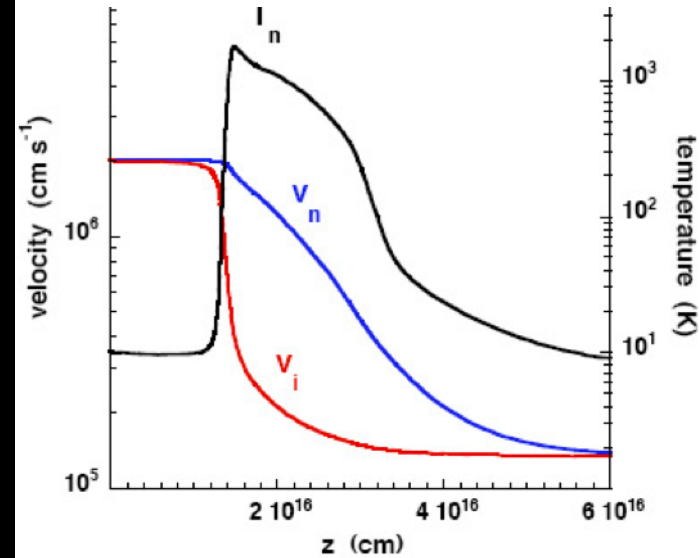
- $V > V_{\text{magnet}}(B)$

J (Jump)-shock = monofluid
Fast dissipation by n-n collisions
High $T_{\text{max}} \sim V^2$



- $V < V_{\text{crit}}(B)$

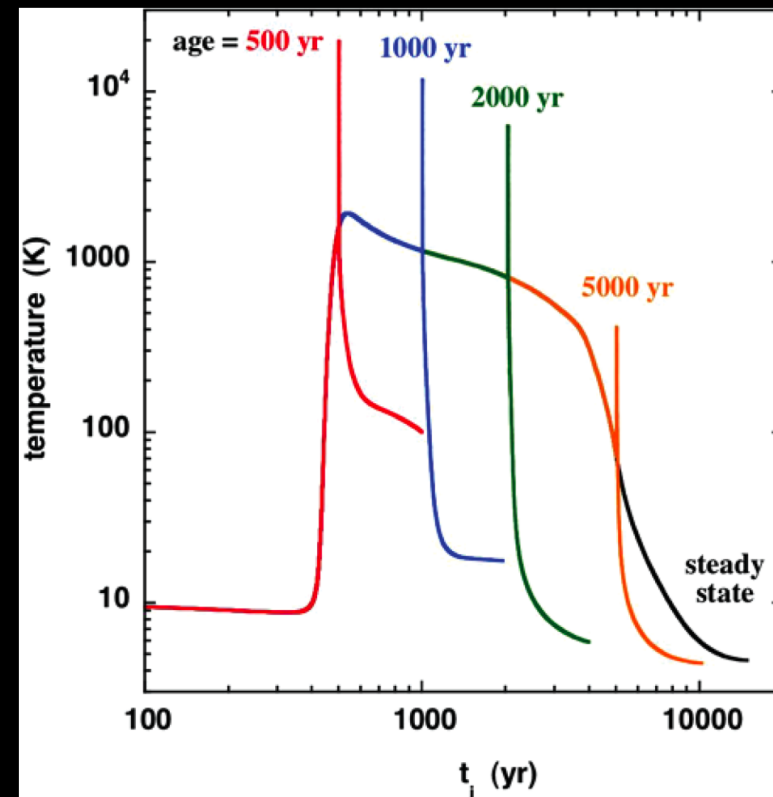
C (Continuous)-shock = multifluid
Slow dissipation by n-i collisions
Wider = Lower T_{max} (H_2 cooling)



YOUNG AGE

- When age $< t_{\text{cool}}$:
 - C-type magnetic « precursor » truncated at $t_{\text{ion}} = \text{age}$
 - followed by a J-type front truncated at $t = \text{age}$ (*Chièze et al. 98, Smith 98*)
 - Well described by truncated steady C + J models if nH , V_s constant (*Lesaffre et al. 2004*)

→ Age is free parameter in the Paris-Durham model



GRAINS

- **Grains** dominate inertia of charged fluid (Guillet 2007)

- For gas/dust ratio ~ 100

$$V_{\text{magnet}} \sim 10 V_{\text{Alfven}} \sim 20 \text{ km/s } \times \mathbf{b}$$

→ **Maximum speed for C-shocks**

- Ion-neutral drift in C-shocks

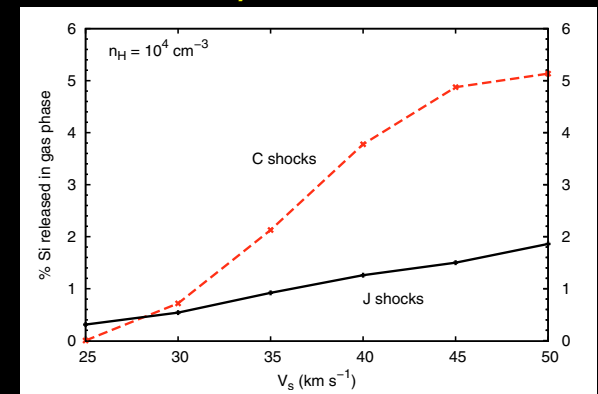
→ Sputtering of ice mantles

- Releases ice species: H_2O , CH_3OH , NH_3 , H_2CO ...

→ Erosion of grain cores

Release Si → SiO (Gusdorf+2008ab)

SiO also released in J-shocks by grain-grain shattering and vaporization (Guillet 2009, 2011)
Not included in public code version



EXTERNAL FUV IRRADIATION

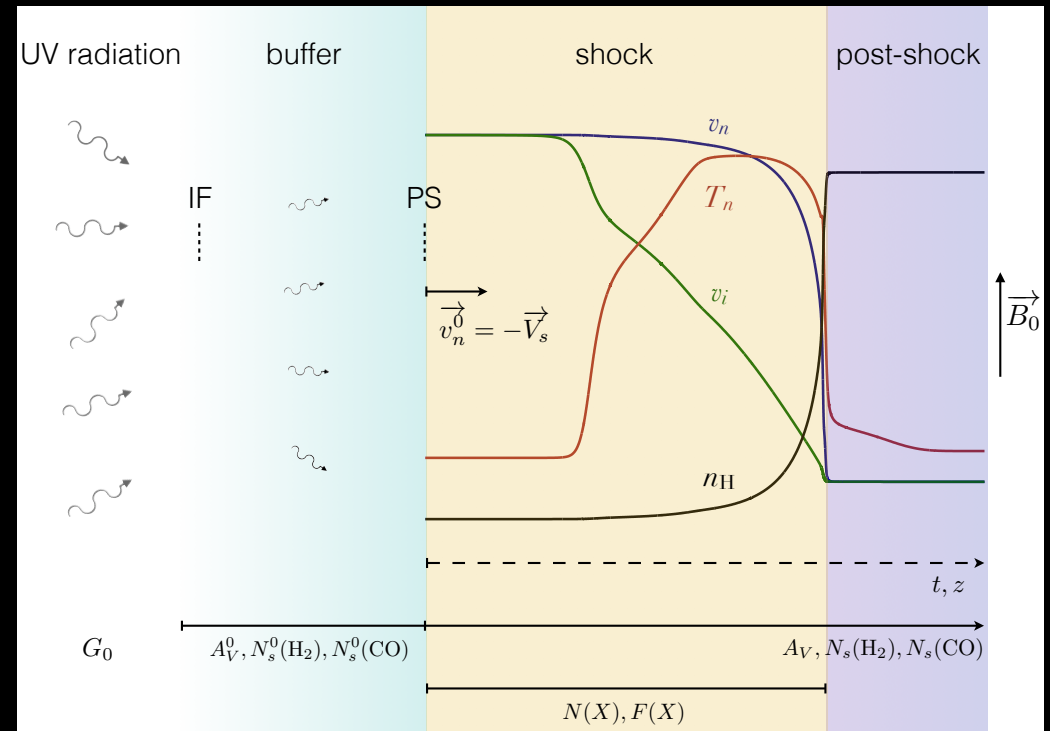
See Godard+2019

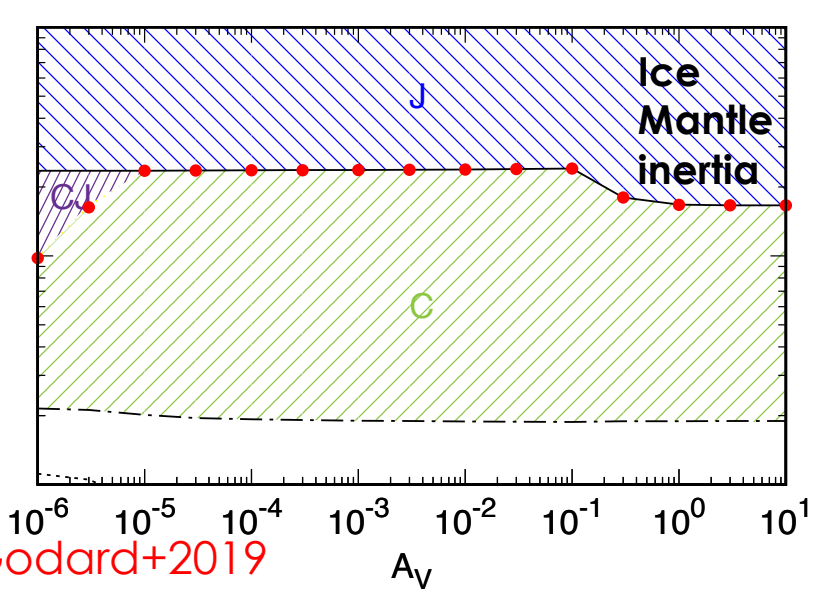
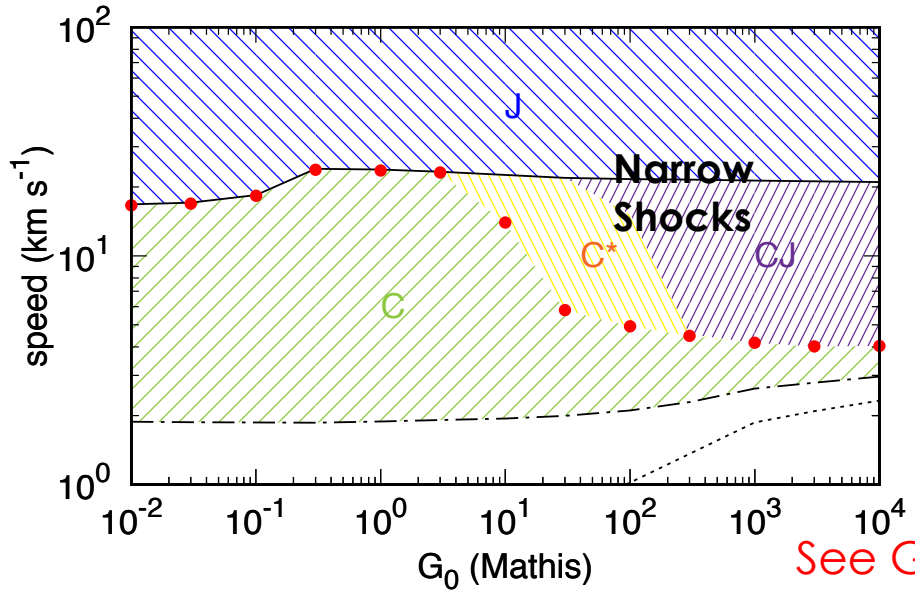
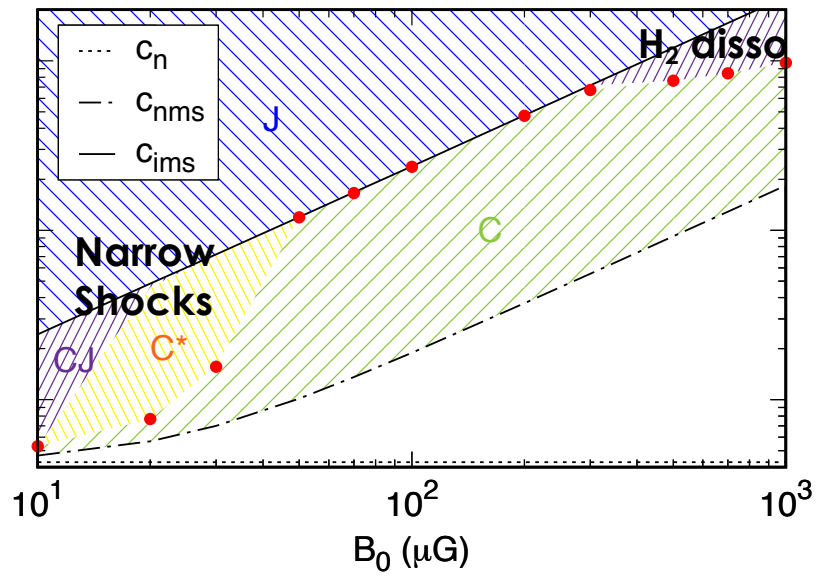
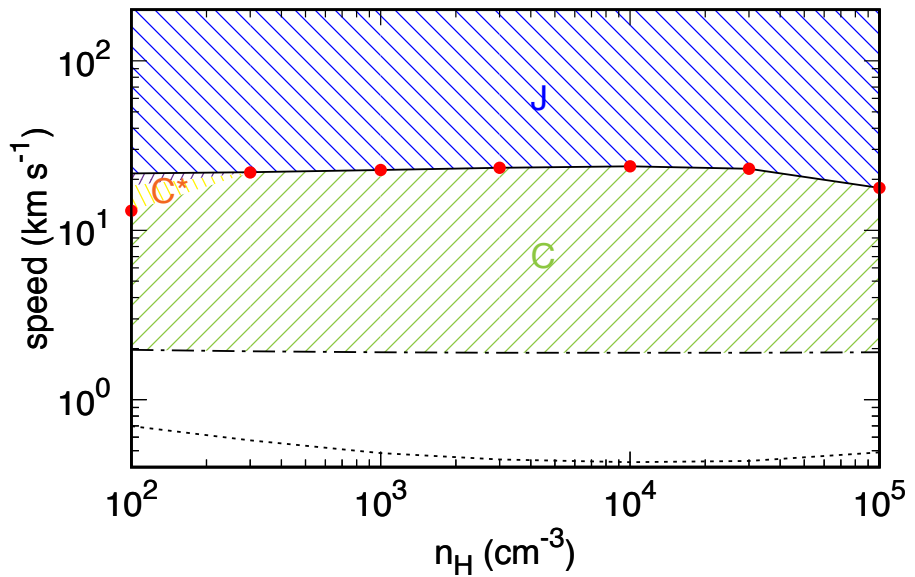
- Extra-Free parameters: G_0 , A_V
- Lowers H₂ cooling + Increases ionization and ion-neutral coupling : narrower hotter C-shock
- → C-shocks disappear for

$$G_0 > 0.2 (n_{\text{H}}/\text{cm}^{-3})^{1/2}$$

and become C* (subsonic) or C-J (jump)

- Specific diagnostics: CH⁺ rotational lines, H₂ rovib (FUV pumping)





See Godard+2019

BACKWARD MODELING (INVERSION PROBLEM)

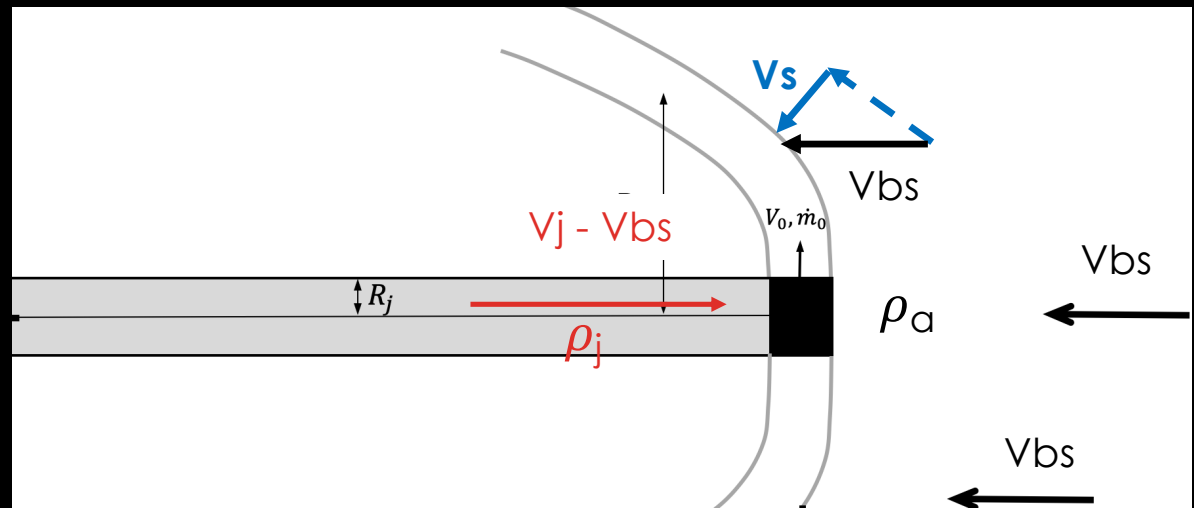
Many free parameters → degeneracy

Ideally combine as many constraints as possible:

- Line fluxes PB: filling factor, depletion in grain+ ice → H₂, CO, O
- Line ratios PB: relative depletion → single specie, or undepleted
- Shock thickness PB: spatial resolution → Hubble, JWST, ALMA
- Line profile width PB: inclination → 2D bowshocks
- Ortho/para of H₂ (if < 3)
- Ambient conditions (nH, G₀, A_v, zeta_CR, zeta_X) from ancillary data
- Shock Age: from speed + size

JET BOWSHOCK

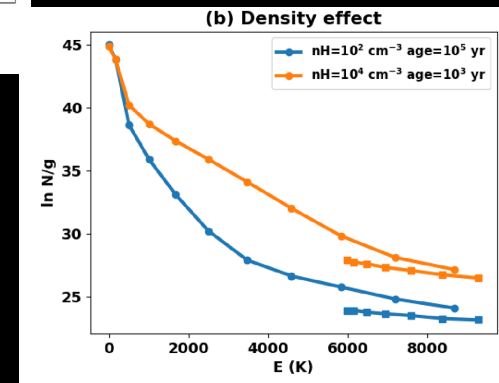
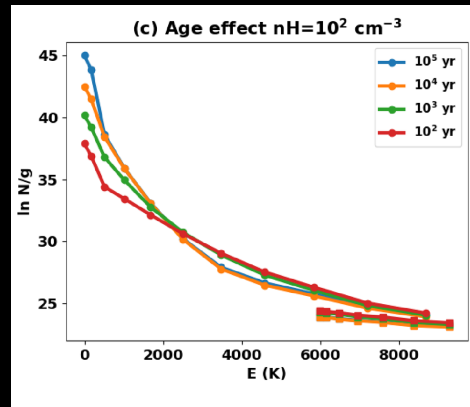
- Lateral escape of shocked jet material
- Swept back by ambient gas
- Oblique shocks : V_s changes along bow surface
- Largest area at low shock speeds



UNRESOLVED 3D BOWSHOCK MODELS

Tram L.N. 2018, MNRAS

- Stitch 1D models along bow surface
- + Compute total H₂ excitation diagram
- resembles slower 1D shock ~ 10 km/s with ~ **same B**
- Fitting with $V_s = V_{\text{bow}}$ would strongly overestimate B-field !
- Distinctive effects of density, age, bow shape, B-field strength and angle → good diagnostic potential

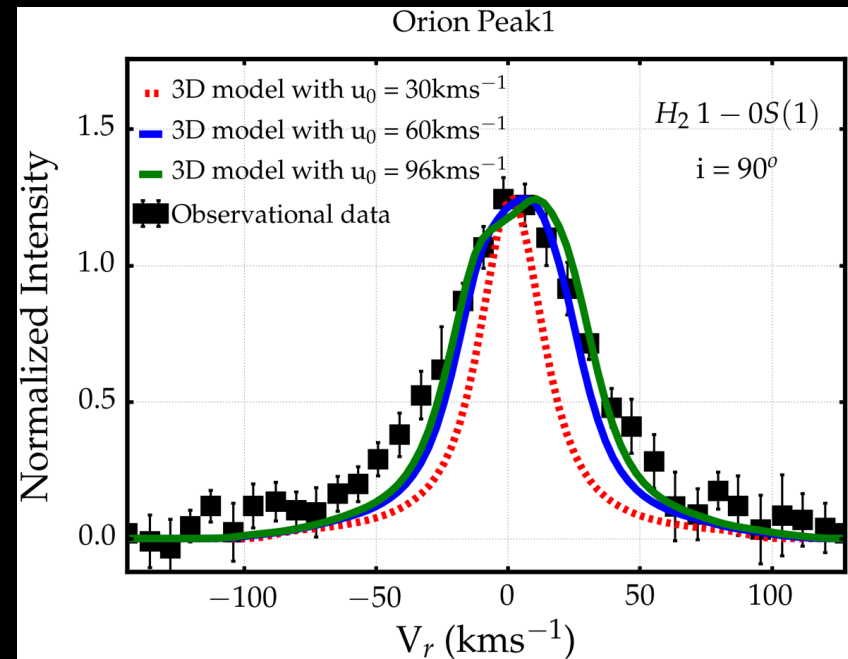
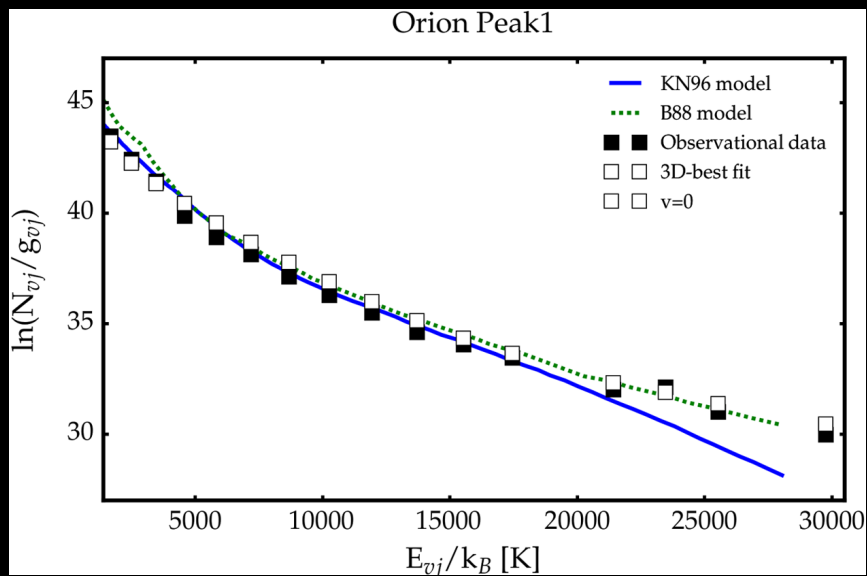


Tram L.N. 2018, MNRAS

FITTING H₂ IN ORION PEAK-1

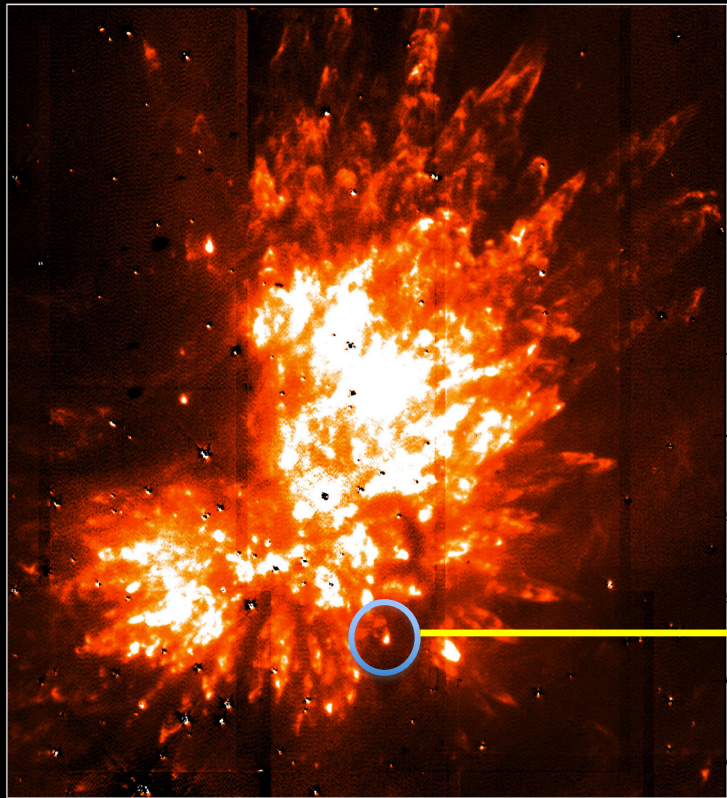
Line profile lifts degeneracy on bow speed U_0
Agrees with proper motion info

- $nH = 10^6 \text{ cm}^{-3}$, age = 1000 yrs, $b = 4.5$
- $U_0 \geq 30 \text{ km/s}$



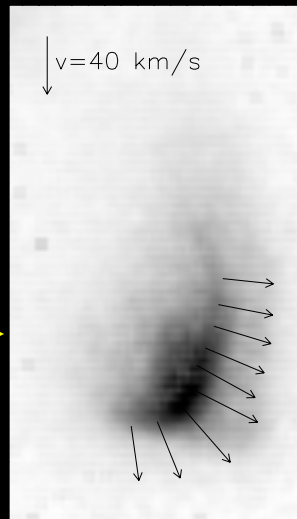
Tram L.N 2018, MNRAS

RESOLVED H₂ BOWSHOCK MODELS



0.15'' resolution
(adaptive optics)

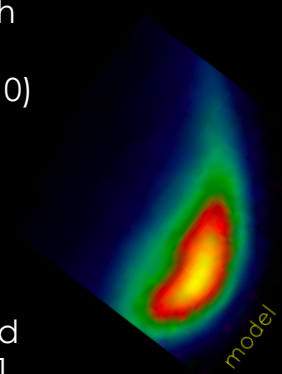
3 Flux-calibrated
H₂ lines ($v=1-0$)



- Resolved emission thickness across bow surface
- Fit cuts by 1D shocks (Kristensen et al. 2008)
- Fit image by 3D bowshock model with projection effects (Gustafsson et al. 2010)

→ C-shocks with
 $b = B/\sqrt{nH} \sim 3-5$

agrees with unresolved
bow model in Peak 1
(and polarization data)



Orion KL

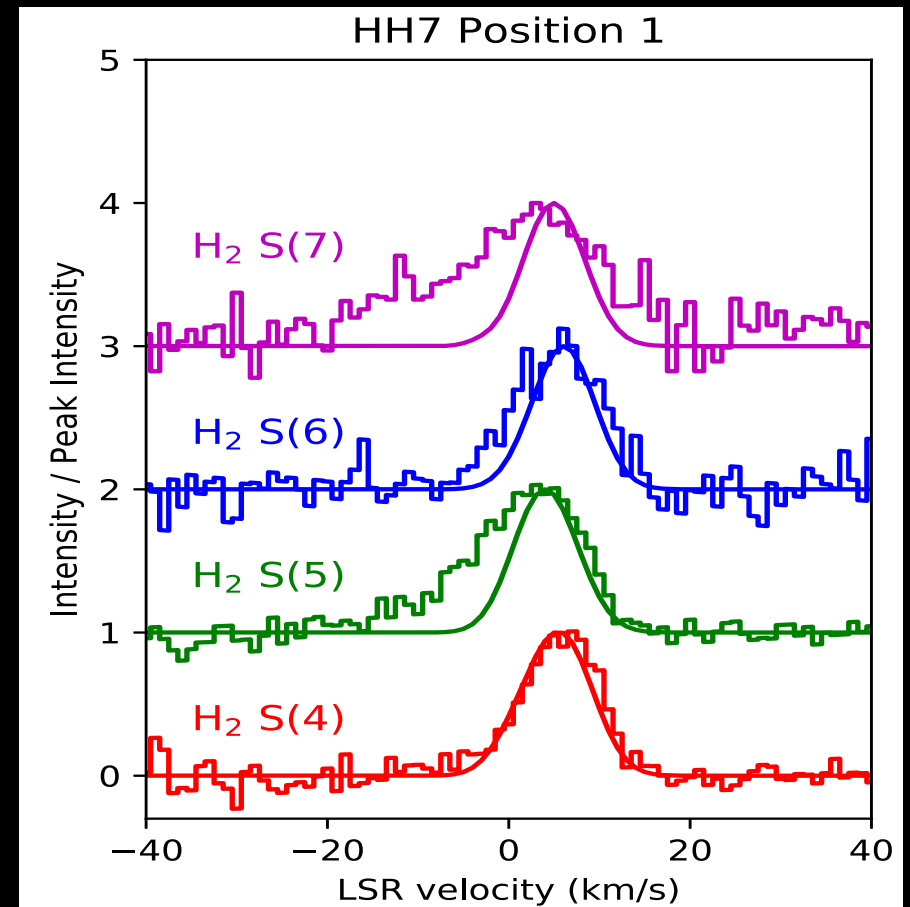
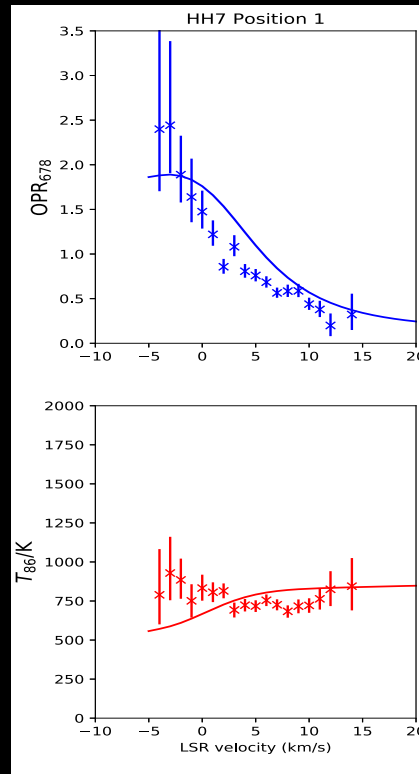
Subaru Telescope, National Astronomical Observatory of Japan

CISCO (H₂ ($v=1-0$ S(1)) – Cont)

January 28, 1999

ORTHO/PARA RATIO OF H_2

- Spectrally resolved observations in HH7 (EXES/SOFIA)
- Velocity shift between ortho and para H_2 lines \rightarrow large OPR variation
- only explained with C-shock model (10^4 cm^{-3} , 20 km/s, 130 microG)
- (Neufeld+2018)

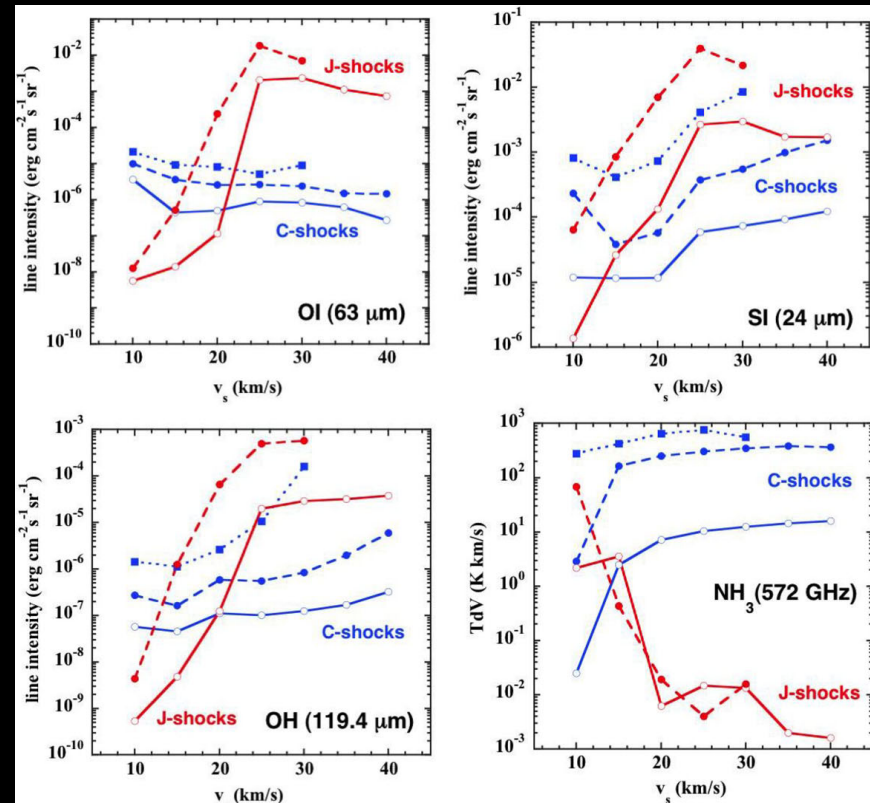


OTHER MOLECULES

- Flower & Pineau des Forêts 2010: CO and H₂O in C and J shocks
- Flower et al. 2010: Methanol in C-shocks
- Flower & Pineau des Forêts 2012: NH₃ emission in C-shocks, **Effect of age, revised methanol and ammonia ice abundances**
- Pineau des Forêts et al. 2013 : OH emission as diagnostic of C vs. J shocks
- Flower et al. 2015 : detailed code description and updated atomic + molecular line predictions using LVG escape formalism

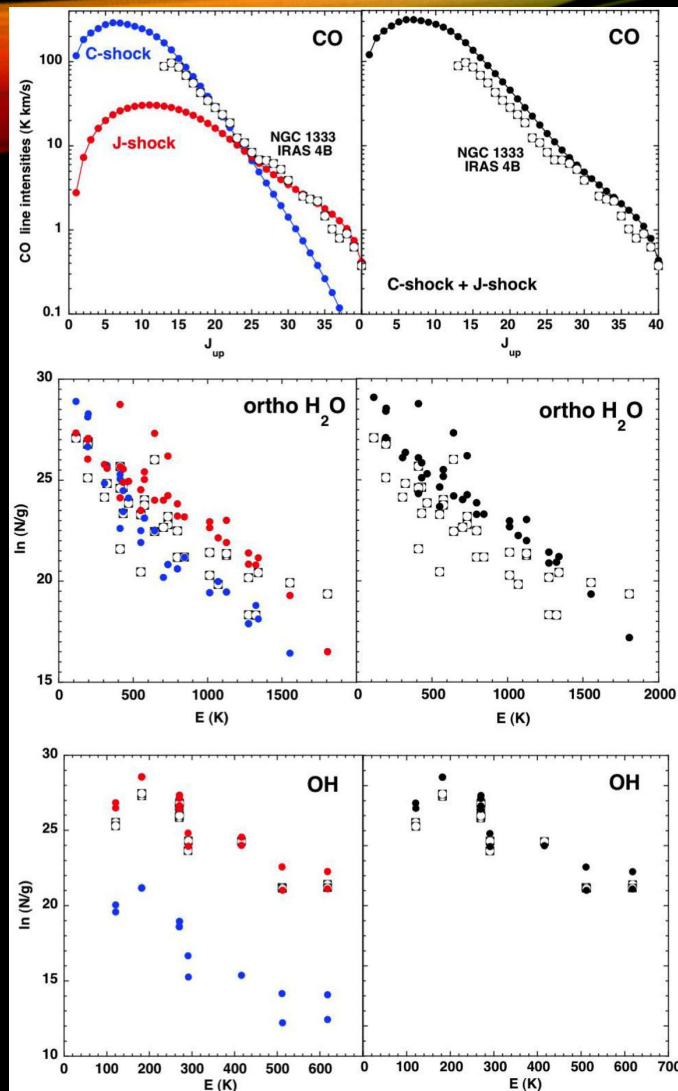
C VS J DIAGNOSTICS

- NB: NH₃ in C-shocks depends on assumed initial ice abundance (ill-known)
- NB2: absolute fluxes depend on « filling-factor » → look at relative line fluxes (excitation diagrams)



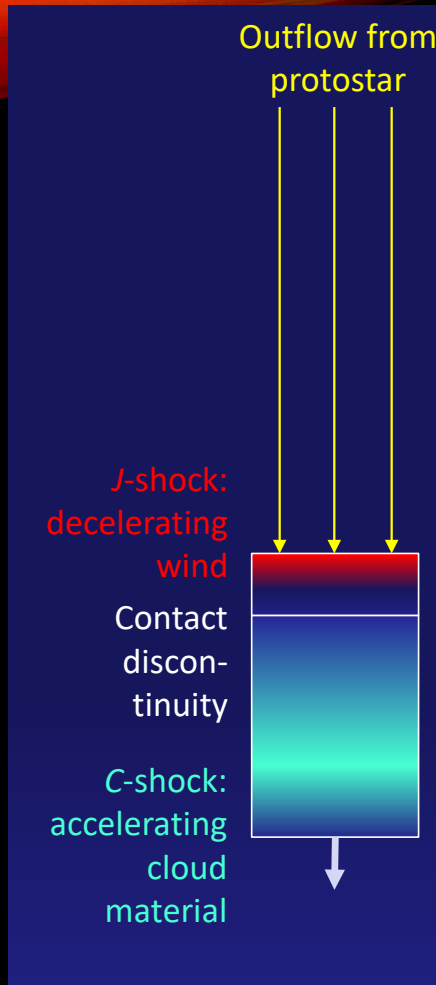
EXAMPLE: IRAS4A

- Herschel PACS data



Flower & Pineau des Forêts 2013

DOUBLE SHOCK STRUCTURE



« BOLOMETRIC METHOD »

% of kinetic energy flux $\frac{1}{2} \rho V_s^3$ radiated by shock tracers:

n_H (cm^{-3})	v_s (km s^{-1})	H ₂ O (per cent)	H ₂ (per cent)	CO (per cent)	O (per cent)	Ly α (per cent)	Total (per cent)
2×10^4	10	1.9	34.2	17.0	0.7	–	53.8
2×10^4	20	11.0	57.0	4.9	0.01	–	72.9
2×10^4	30	6.8	72.7	2.1	–	–	81.6
2×10^4	40	4.3	80.6	1.1	–	–	86.1
2×10^5	10	4.9	26.2	19.6	0.1	–	50.8
2×10^5	20	17.1	48.8	3.9	–	–	69.8
2×10^5	30	9.6	69.3	1.4	–	–	80.4
2×10^5	40	5.9	79.5	0.7	–	–	86.1
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2×10^4	10	3.2	75.5	6.8	–	–	85.5
2×10^4	20	15.0	72.5	9.0	0.01	–	96.6
2×10^4	30	18.7	3.3	0.8	6.6	41.1	80.4
2×10^5	10	5.7	64.5	3.6	–	–	73.8
2×10^5	20	50.7	15.1	5.8	1.8	–	73.3
2×10^5	30	23.3	3.0	0.1	2.6	48.2	77.3

- H₂ main coolant in non-dissociative shocks (25%-80%): useful « bolometric » tool at $V_s < 30 \text{ km/s}$

- CO % and H₂O % depend too strongly on (n_H, V_s) (but improve total % if added to H₂)

- O and H lines → dissociative J-shocks

C-type

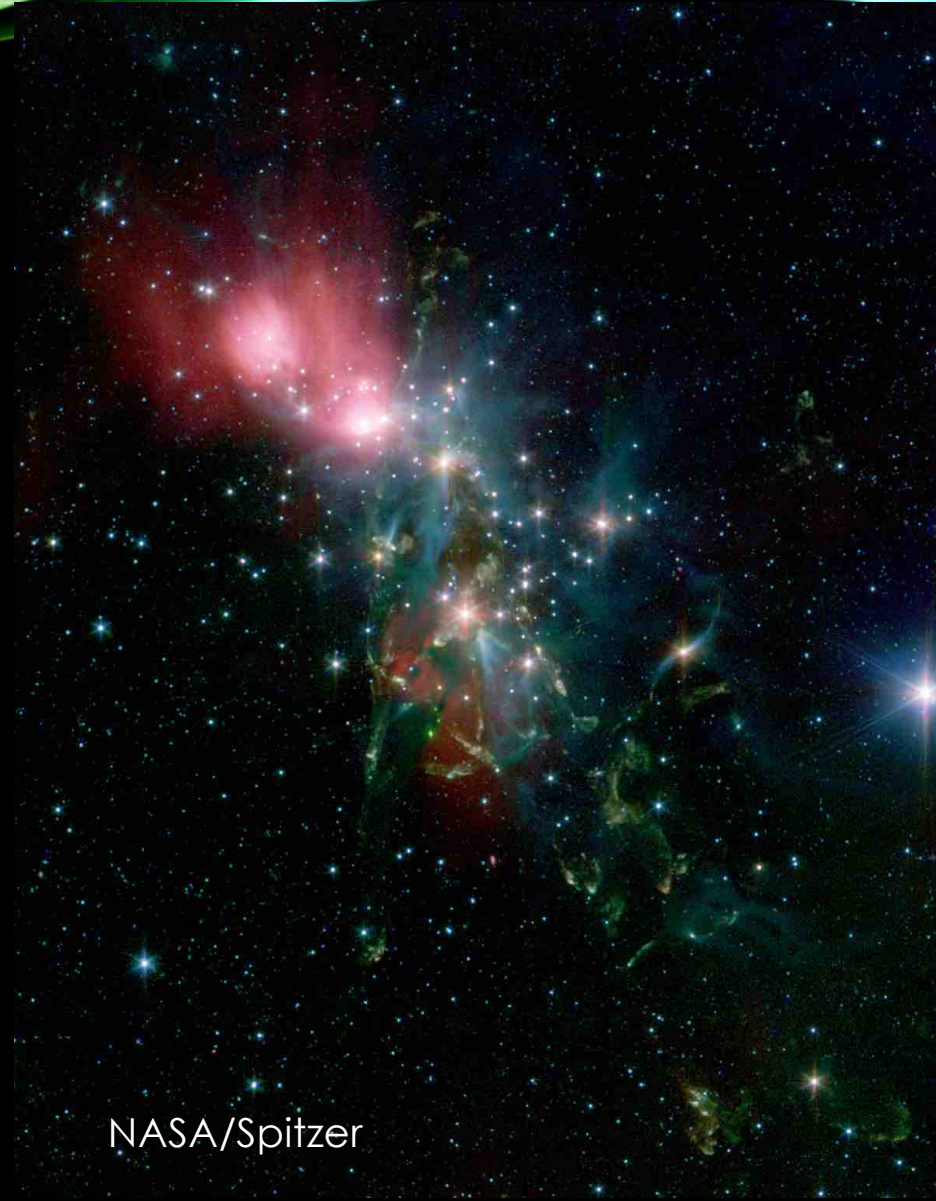
J-type

H₂ BOLOMETRIC METHOD APPLIED

Maret et al. 2009

- *Spitzer*: H₂ pure rotational lines from outflows in NGC1333 (green in image)
C-shocks at $V_s \sim 20\text{-}30$ km/s
- $L(\text{H}_2) \rightarrow \text{KE flux in shock} = L_{\text{mech}}$
 \rightarrow Momentum injection rate :
$$P_s = 2 L_{\text{mech}} / V_s$$

$$= 2\text{-}5 \cdot 10^{-5} M_{\odot} \text{ km/s /yr per flow}$$
- **After 2×10^5 yr, each flow imparts 1 km/s turbulence in $5 M_{\odot}$**
20 flows: $100 M_{\odot} = 30\%$ of cloud mass

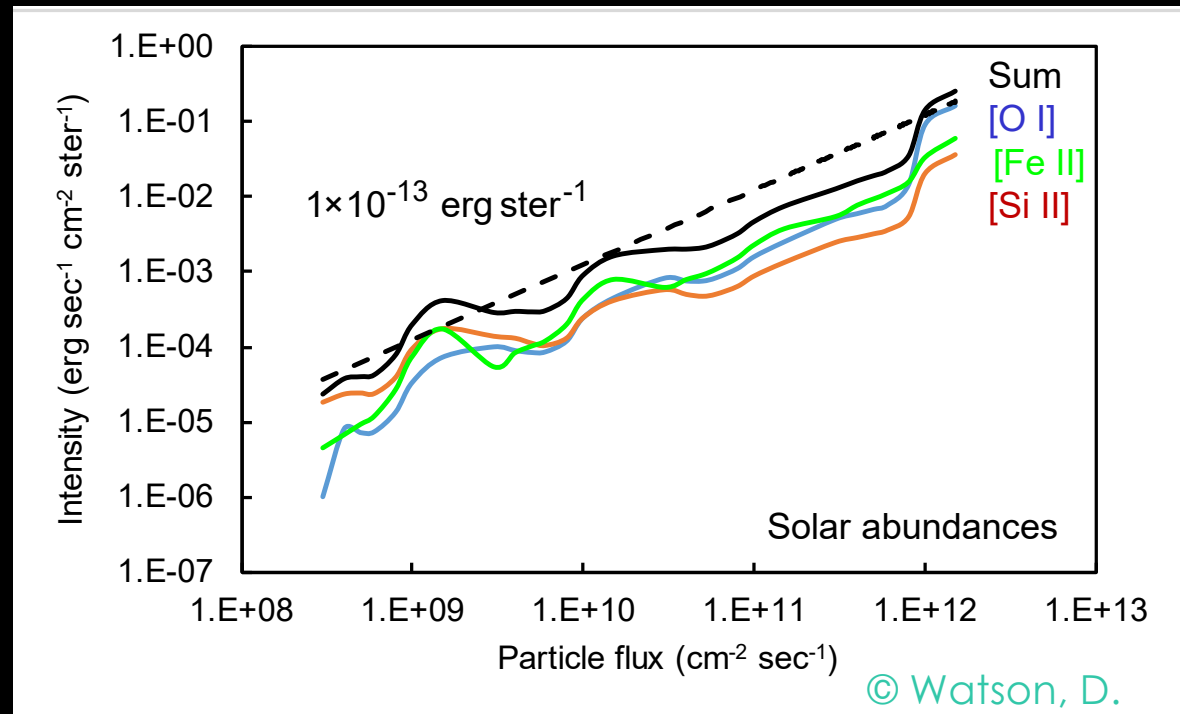


NASA/Spitzer

« BOLOMETRIC METHOD » FOR DISSOCIATIVE SHOCKS

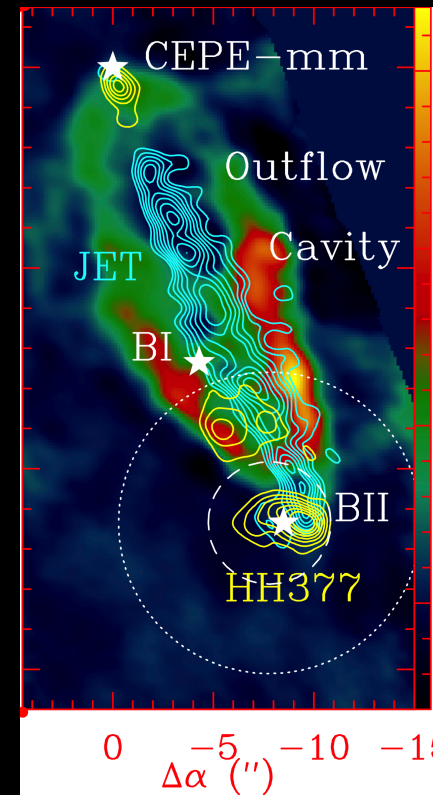
- [O I] 63mic flux is (roughly) proportional to **mass-flux into shock** (Hollenbach 1985)
- Same for [Fe II]26mic and [Si II]35mic

$$\begin{aligned} \dot{M}_w &= 8.1 \times 10^{-5} L ([O I] 63.2 \mu\text{m}) \frac{M_\odot \text{ year}^{-1}}{L_\odot} \\ &= 9.8 \times 10^{-4} L ([Si II] 34.8 \mu\text{m}) \frac{M_\odot \text{ year}^{-1}}{L_\odot} \\ &= 1.4 \times 10^{-3} L ([Fe II] 26.0 \mu\text{m}) \frac{M_\odot \text{ year}^{-1}}{L_\odot}. \end{aligned}$$

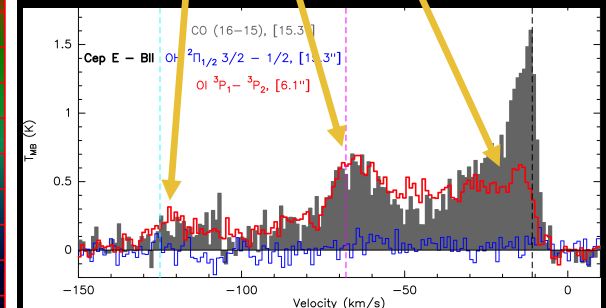


LIMITATIONS OF « BOLOMETRIC » METHODS

- Several shock fronts contribute to [O I] not just the final terminal shock
 - Overestimate mass-flux in jet/wind
- Need spectrally (SOFIA-GREAT, EXES) and spatially (JWST, ALMA, NOEMA) resolved observations of shocks to separate components



- 3 components in [O I] profile (SOFIA)
Gusdorf+2017



CONCLUSIONS

- Interstellar shock waves are complex systems to model, but very detailed public 1D models exist for different shock regimes
- Backward inversion is difficult exercise
 - Many free parameters: beware of degeneracies, combine constraints !
 - Several key atomic coolants (Fe+, S, Si+) may have grain depletion
 - When non-dissociative: Output radiation depends on preshock chemical conditions (ice mantles + ambient G0 and zeta): necessary to use self-consistent preshock before running models
 - May have several shock fronts in beam (reverse J + forward C): favor high angular resolution and spectral resolution: JWST, SOFIA and ALMA
- Fast progress in 2D/3D simulations including NEQ chemistry: new tool